

# Quantum Mechanics

Dual nature of Matter waves

Davisson &

Germer experiment

Matter waves & Uncertainty

photo electric effect

gamma ray microscope

principle

Heisenbergs uncertainty principle

matter wave properties

phase velocity & group velocity  
expressing

Eigen values & Eigen functions

$\hat{H}$

properties of wave functions

Schrodinger wave equation

Time independent

Time dependent

Time dependent

Compton effect

Inadequacy of Classical physics

photo electric effect

Einstein PEE

Inadequate = not enough



SAGS

- ① Difference bw phase velocity and group velocity
- ② postulates of Quantum mechanics
- ③ Derive Einsteins photo electric Equation
- ④ state and Explain Heisenbers uncertainty principle.
- ⑤ physical significance of wave function
- ⑥ properties of Matter waves
- ⑦ Heisenb. Gamma ray microscope
- ⑧ Relation bw phase velocity and group velocity.

# Quantum Mechanics

**Paper**

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3 parts

part - 1 (1) PEE & (2) Compton effect (3) Ein  
(3) Einstein PEE  $10m$   
(4) Quantum mechanics postulates  $10m$

part - 2 :-

w (1) Debroglie wavelength  
w (2) Heisenberg uncertainty principle  
(3) matter waves properties  $10m$   
(4) LA phase velocity & group v difference  $3-5m$   
w (5) Davison germe experiment  
(6) Gamma ray microscope.  $10m$  ✓

part - 3 :- w (1) Schrony time d

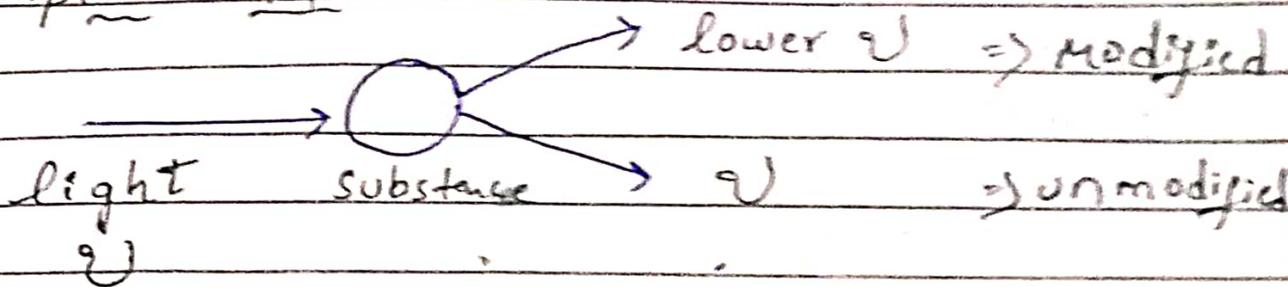
w (2) Time In

3) LA properties of wave function

4) LA eigen values eigen functions.

① what is Compton Effect? Considering Compton effect/scattering, derive an expression for shift.

Compton Effect :-

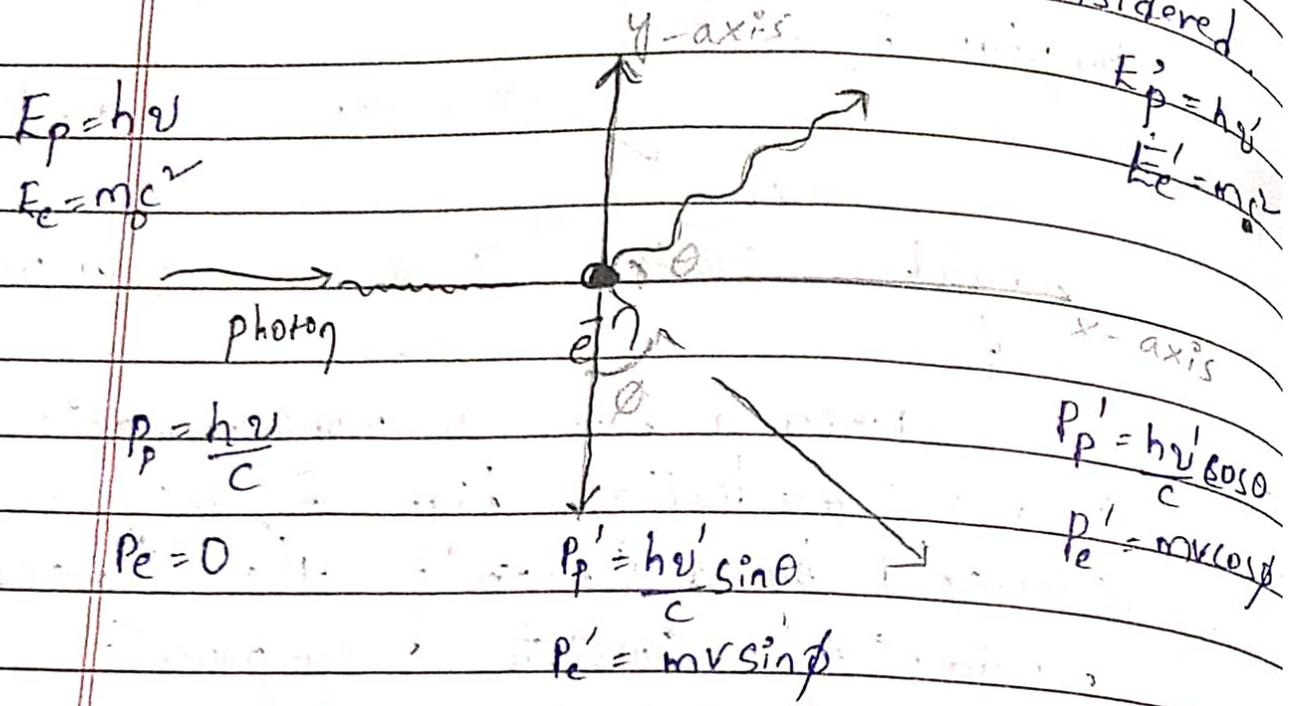


When a beam of monochromatic source radiation high frequency is scattered by a substance, the scattered radiation consists 2 components

- (i) one having lower frequency or greater wavelength which is called as Modified radiation
- (ii) The other having same frequency as original (or) same wavelength it is called unmodified radiation.

This phenomenon is called Compton effect.

Compton assumed that the  $e^-$  is free and at rest before with photon. After collision the relative mass of  $e^-$  is considered



Let a photon of energy  $h\nu$  collides with an  $e^-$  at rest.

During collision some part of energy of photon is given to  $e^-$ , because of which  $e^-$  gain KE and recoils.

The process of recoiling of  $e^-$  and scattering of photon is shown in fig.

$\theta$  is scattering angle of photon &  
 $\phi$  is recoil angle of  $e^-$ .

[Recoil =  $\vec{p}_p - \vec{p}_p'$  वा  $\vec{p}_e'$  वेक्टरों की तरह reverse]

$$m_1 v_1 \cos \phi = m_1 v_1 - m_1 v_1 \cos \theta$$

and (4) solve and adding

$$\sin^2 \phi + m^2 v^2 \cos^2 \phi = h^2 v^2 \sin^2 \theta = h^2 v^2$$

$$2(\sin^2 \phi + \cos^2 \phi) = h^2 v^2 (\sin^2 \theta + \cos^2 \theta)$$

$$2c^2 = h^2 v^2 + h^2 v^2 - 2h^2 v_1 v^1 \cos \theta$$

$$2c^2 = h^2 [v^2 + v_1^2 - 2v v_1 \cos \theta]$$

of the system Before collision =  $h v_1 + m_1 v_1^2$

of the system After collision =  $h v + m v^2$

According to law of conservation of Energy

Case (i) :- when  $\theta = 0$  ( $\cos \theta = 1$ )  $\Rightarrow$

$$\Delta \lambda = \frac{h}{m_0 c} (1 - \cos \theta) \Rightarrow \Delta \lambda = \frac{h}{m_0 c} (1 - 1) = 0$$

$$\Delta \lambda = 0 \text{ at } \theta = 0$$

$\therefore$  no change in frequency/wavelength of incident radiation.

$\therefore$  no scattering along direction of incidence.

Case (ii)

$$\theta = \frac{\pi}{2} \Rightarrow \cos \theta = 0$$

$$\Delta \lambda = \frac{h}{m_0 c} (1 - 0) = \frac{h}{m_0 c}$$

$$\Delta \lambda = \frac{6.624 \times 10^{-34}}{9.1 \times 10^{-31} \times 3 \times 10^8}$$

$$\Delta \lambda = 0.2426 \times 10^{-10} \text{ m}$$

$$\Delta \lambda = 0.02426 \text{ \AA}$$

Compton wavelength

Case (ii)  $\theta = \pi$

$$\cos \theta = -1$$

$$\Delta \lambda = \frac{h}{m_0 c} (1 - \cos \theta)$$

$$= \frac{h}{m_0 c} (1 - (-1))$$

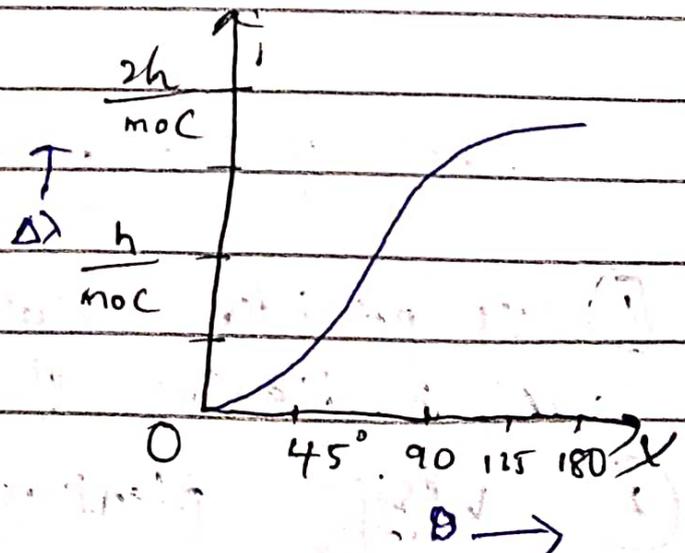
$$\Delta \lambda = 2 \cdot \frac{h}{m_0 c}$$

$$\Delta \lambda = 0.04852 \text{ \AA}$$

when  $\theta$  change from  $0$  to  $2\pi$

$\lambda$  also changes

$$\lambda = n \frac{h}{m_0 c}$$



frequency of Scattered photon:

from  $\frac{1}{\nu'} - \frac{1}{\nu} = \frac{h}{m_0 c^2} (1 - \cos \theta)$

$\frac{1}{\nu'} = \frac{1}{\nu} + \frac{2h \nu \sin^2 \theta / 2}{m_0 c^2}$

$\frac{1}{\nu'} = \frac{m_0 c^2 + 2h \nu \sin^2 \theta / 2}{m_0 c^2 \nu}$

$\nu' = \frac{m_0 c^2 \nu}{m_0 c^2 + 2h \nu \sin^2 \theta / 2}$

$\therefore \nu' = \frac{\nu}{1 + \frac{2h \nu \sin^2 \theta / 2}{m_0 c^2}}$

$\nu' = \frac{\nu}{1 + 2d \sin^2 \theta / 2}$        $d = \frac{h \nu}{m_0 c^2}$

Uses of Compton effect:

- ① It provides evidence of electromagnetic radiation of particle.
- ② Verifies plankus Quantum hypothesis.

③ verifies relations 'Indepently'

$$E = mc^2 \quad \text{and} \quad m = \frac{m_0}{\sqrt{1 - \frac{v^2}{c^2}}}$$

② Derive Einstein's photo electric equation  
According Einstein's explanation in photo electric effect, one photon is completely absorbed by one  $e^-$  which there by gains quantum energy and may be emitted from metal surface. The photon energy split into 2 parts.

(i) A part of energy is used to separate  $e^-$  from metal surface.

This energy is called work function  $W_0$

(ii) Another part of energy is used to give KE to  $e^-$

$$E \text{ of photon} = h\nu$$

$$h\nu = W_0 + KE \quad \text{--- (1)}$$

when In some case where photon have an amount of energy value where it can only emit  $e^-$  from metal

Then KE of  $e^-$  will be zero

$$\frac{1}{2}mv^2 = 0 \Rightarrow h\nu_0 = W_0 \quad [\nu = \nu_0]$$

②

Sub (2) in (1)

$$h\nu = h\nu_0 + \frac{1}{2}mv^2$$

$$\boxed{\frac{1}{2}mv^2 = h\nu - h\nu_0 = h(\nu - \nu_0)}$$

from (1)  $\frac{1}{2}mv^2 = h\nu - h\nu_0$  - (4) for a particular metal work function  $h\nu_0$  is constant

$$\Rightarrow \frac{1}{2}mv^2 \propto h\nu \Rightarrow v^2 \propto \nu$$

$\therefore$  velocity of photo electron is proportional to frequency of incident radiation.

If  $V_0$  be stopping potential then

$$\frac{1}{2}mv^2 = eV_0 - (5)$$

from (3) & (5)  $eV_0 = h\nu - h\nu_0$

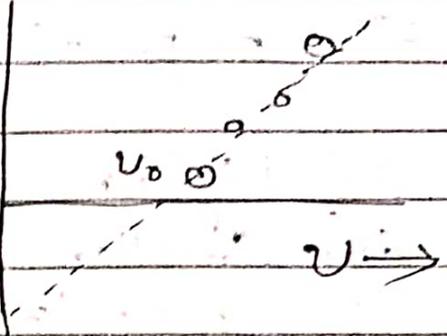
$$\boxed{V_0 = \frac{h\nu}{e} - \frac{h\nu_0}{e}} - (6)$$

eqn (1) (3) (6) are different forms of Einstein's P.E.E.

from (1)  $V_0 = \frac{h}{e} \nu - \frac{h}{e} \nu_0$

it represents a straight line.

↑  
 $V_0$



(3)

postulates of Quantum Mechanics.  
The physics is study of nature & the study which we did till 1900's called classical physics.

→ But classical physics had some rules which failed to prove atomic substance in detail.

On Dec 1900 max planck gave his presentation where we can study details of atoms on black body radiation & leading to Quantum mechanics.

According to planck's :-

(1) Emission of radiation is due to vibration of charged particles in body.

(2) The radiator body contains SHO of all frequencies.

(3) The particles contains tiny packet of energies called as photons / Quantum

(4) Energy of photon is  $E = h\nu$

(5) general for oscillator  $E = nh\nu$   $n = 0, 1, 2, 3, \dots$

(6) later Albert Einstein explained PEE

Bohr introduced Bohr atomic model

compton explained Compton effect

Based on planck's theory & Quantum mechanics of planck

## (4) properties of matter waves

- (i) Lighter particle have greater wavelength
- (ii) - Smaller velocity particle have greater wavelength
- (iii) Velocity = 0 then such particle have infinity wavelength
- (iv)  $v = \alpha \Rightarrow$  wavelength = 0  $\checkmark$  of a particle  $\Rightarrow$  wave is generated by motion of particle
- (v) Velocity of matter wave depends on  $v$
- (vi) Velocity of matter particle.
- (vii) Velocity of matter wave is greater than velocity of light.
- (viii) A particle in motion with associated matter wave has two different velocities.

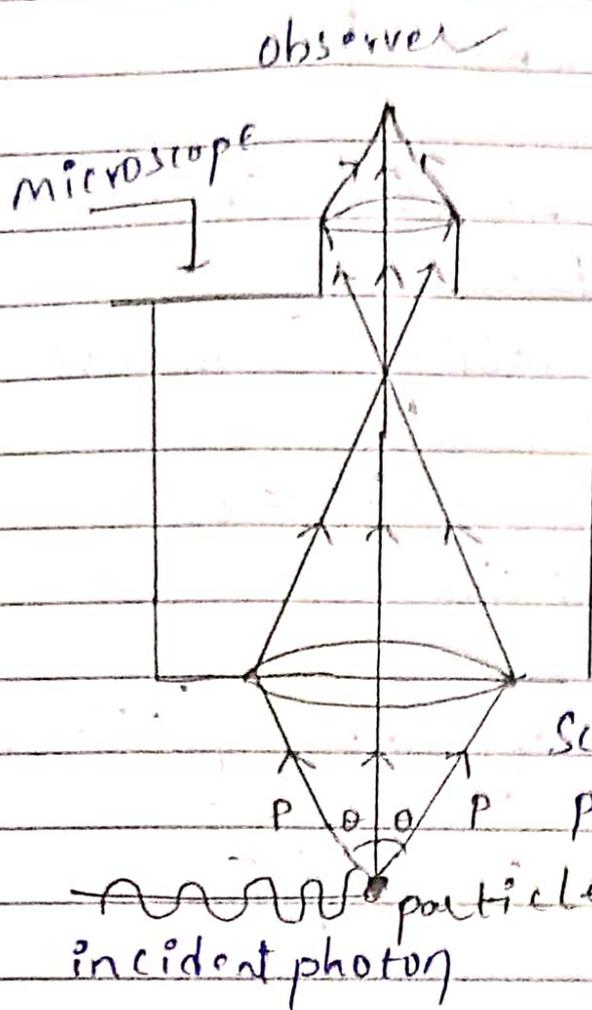
(i) mechanical motion of particle ( $v$ ) (ii) wave velocity ( $w$ )  
we know

$$E = h\nu \quad \& \quad E = mc^2 \Rightarrow h\nu = mc^2 \Rightarrow \boxed{\nu = \frac{mc^2}{h}}$$

$$\text{but } w = \nu \lambda \Rightarrow w = \frac{mc^2}{h} \times \frac{h}{m v}$$

$$\boxed{w = \frac{c^2}{v}}$$

⑤ write about Gamma ray microscope



The position and linear momentum of an  $e^-$  can be measured using microscope with angular aperture  $2\theta$  as showing fig.

Resolving power of microscope is

$$\Delta x = \frac{\lambda}{2 \sin \theta} \quad (1)$$

$e^-$  may be observed when a photon is scattered by it into field of view of microscopic lens.

when a photon of initial momentum  $p = \frac{h}{\lambda}$ , after scattering enters the field of view of microscope

It may be anywhere within angle  $2\theta$ . Thus,  $p_n$  may lie b/w  $p \sin \theta$  and  $-p \sin \theta$

According to law of Conservation of momentum, the uncertainty in momentum is given by.

$$\Delta p_n = p \sin \theta - (-p \sin \theta)$$

$$\Delta p_n = 2p \sin \theta$$

$$\Delta p_n = \frac{2h}{\lambda} \sin \theta \quad \left[ \because \lambda = \frac{h}{p} \right]$$

from (1) & (2) we get

$$\Delta n \cdot \Delta p_n \approx \frac{1}{2 \sin \theta} \times \frac{2h}{\lambda} \sin \theta$$

$$\Delta n \cdot \Delta p_n \approx h$$

Q6) what are phase velocity and group velocity obtain their expressions for them

(a) phase velocity / wave velocity ( $v_p$ )

The velocity with which planes of constant phase advances through medium is known as phase velocity.  
 (or)

The ratio of angular frequency and propagation constant is called phase velocity, it denoted by  $v_p$

Let us consider a wave of single-frequency and wavelength, travels through a medium then displacement of wave is given by

$$y = a \sin(\omega t - kn)$$

for planes of constant phase,

$$\omega t - kn = \text{const}$$

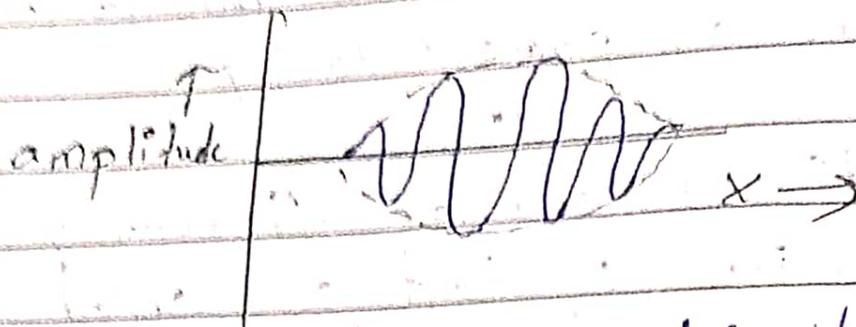
on differentiating wrt time,

$$\omega \frac{dt}{dt} - k \frac{dn}{dt} = 0$$

$$\frac{\omega}{k} = \frac{dn}{dt} \Rightarrow \boxed{\frac{dn}{dt} = \frac{\omega}{k}} \quad \boxed{v_p = \frac{\omega}{k}}$$

(ii) Group velocity is  $V_g$

The velocity with which the energy in group is transmitted is known as group velocity. It is denoted by  $V_g$ .



Schrodinger explained dual nature of matter in terms of wave packet

A wave packet contains a group of waves slightly differing in velocity and wavelength, with phases and amplitudes chosen in such a way that they interfere constructively over a small region of space where the particle can be located and outside this space they interfere destructively so that amplitude reduces to zero. Such a wave packet moves with its own velocity called group velocity.

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The individual waves forming the wave packet have an average velocity called phase velocity:

Consider two waves having same amplitude 'a' slightly different angular frequencies ( $\omega_1$  and  $\omega_2$ ) and phase velocities ( $v_1$  and  $v_2$ ). Then the waves can be represented as

$$y_1 = a \sin(\omega_1 t - k_1 x) \quad \text{--- (1)}$$

$$y_2 = a \sin(\omega_2 t - k_2 x)$$

$\therefore$  According to superposition principle resultant displacement (y)

$$y = y_1 + y_2$$

$$y = a \sin(\omega_1 t - k_1 x) + a \sin(\omega_2 t - k_2 x)$$

$$y = 2a \cos \left[ \left( \frac{\omega_1 + \omega_2}{2} \right) t - \left( \frac{k_1 + k_2}{2} \right) x \right]$$

$$\left( \sin C + \sin D = 2 \cos \frac{C-D}{2} \sin \frac{C+D}{2} \right) \sin \left[ \left( \frac{\omega_1 + \omega_2}{2} \right) t - \left( \frac{k_1 + k_2}{2} \right) x \right]$$

$$y = 2a \cos \left( \frac{d\omega}{2} t - \left[ \frac{dk}{2} \right] x \right) \sin(\omega t - kx)$$

$\therefore$  Amplitude =  $2a \cos \left( \frac{d\omega}{2} t - \frac{dkx}{2} \right)$

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$$\text{Amplitude} = 2a \cos\left(\frac{d\omega}{2} \left(t - \frac{dk}{2} x\right)\right)$$

$$\text{Amplitude} = 2a \cos \frac{d\omega}{2} \left(t - \frac{dk}{2} x\right)$$

$$\text{Amplitude} = 2a \cos \frac{d\omega}{2} \left(t - \frac{x}{v_g}\right)$$

$$v_g = \frac{d\omega}{dk} = \frac{\omega_1 - \omega_2}{k_1 - k_2}$$

Relation. bw. group velocity and phase velocity :-  
we know,

$$v_p = \frac{\omega}{k}$$

$$\Rightarrow \omega = v_p k \Rightarrow d\omega = v_p dk + k dv_p$$

$$\frac{d\omega}{dk} = v_p + k \frac{dv_p}{dk}$$

$$v_g = v_p + k \frac{dv_p}{dk} \quad \text{--- (1)}$$

$$\frac{d}{dx} \frac{1}{x} = -\frac{1}{x^2}$$

we know  $k = \frac{2\pi}{\lambda} \Rightarrow \lambda = \frac{2\pi}{k} \Rightarrow d\lambda = -\frac{2\pi}{k^2} dk$

$$\frac{d\lambda}{dk} = -\frac{2\pi}{k^2} \quad \text{--- (2)}$$

from (1)  $v_g = v_p + k \frac{dv_p}{d\lambda} \frac{d\lambda}{dk}$  --- (2)

from (2) & (3)

$$V_g = V_p + K \cdot \frac{dV_p}{d\lambda} \left( -\frac{2\pi}{K\lambda} \right)$$

$$V_g = V_p - \frac{dV_p}{d\lambda} \frac{2\pi}{K} \quad K = 2\pi$$

$$V_g = V_p - \lambda \frac{dV_p}{d\lambda} \quad \checkmark$$

$$K = 2\pi \Rightarrow \lambda = \frac{2\pi}{K}$$

$$V_g = V_p + K \cdot \frac{dV_p}{dK}$$

① write about properties of wave-function and physical significance of wave function

→ Wave function is a mathematical function introduced by Schrodinger to explain debroglies Ideas of wave.

$$\psi(x, t) = A e^{-i(kx - \omega t)}$$

properties of wave function:-

Wave function  $\psi(x, y, z, t)$  is

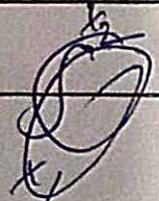
The probability that a particle will be found <sup>at a place</sup> in space at a given time

(1) wave function must be finite or zero at any point

(2) wave function must have a single valued

(3) wave function be a continuous wave. function poses normalize condition

$$\int_{-\infty}^{\infty} |\psi|^2 dx dy dz = 1$$

→ The probability of find system b/w  $x_1$  &  $x_2$  is 

$$P = \int_{n_1}^{n_2} |\psi|^2 dn$$

(b) probability per unit volume is probability density.

$$\Rightarrow P = |\psi|^2 dv$$

$$\int P dv = 1$$

$$\int |\psi|^2 dv = 1$$

physical significance of wave function

→ By using wave function we can get useful information of system

probability density:-

The wave function is in itself complex

$\psi = A + iB$  so its complex conjugate

$$\psi^* = A - iB$$

$$|\psi|^2 = \psi \psi^*$$

$$= (A + iB)(A - iB)$$

$$= A^2 + B^2 \text{ Real quantity}$$

$|\psi|^2$  probability density

The probability of finding a particle at any point at a given time is proportional to probability density

$$|\psi|^2 = \psi\psi^*$$

A large value  $|\psi|^2$  means strong possibility of finding particle at the given point of time.

→ The wave function  $\psi$  is different from zero in a finite region which is called wave packet

② Explain Eigen values and Eigen functions also Expectation values.

There is a class of functions  $\psi$  which are called Eigen functions.

If for an operator  $\hat{O}$  and a constant  $\lambda$

$$\hat{O} \psi(x) = \lambda \psi(x)$$

then,  $\lambda$  is called eigen value of operator  $\hat{O}$   
 $\psi(x)$  is eigen function of  $\hat{O}$

In Schrodinger wave equation in operator form, we write,

$$\hat{H} \psi = E \psi$$

where  $\hat{H} = \frac{-\hbar^2}{2m} \nabla^2 + V$  and

$$E = i\hbar \frac{\partial}{\partial t}$$

So we can define eigen values & eigen functions as

$$\hat{H}\psi = E\psi$$

↓ eigenvalue  
→ eigen-function

The values of Energy  $E_n$  which can solve schrodinger equation is called eigen value and corresponding function is called eigen value-function

# Davisson-Germer Experiment

[Experimental Confirmation of matter wave]

→ Davisson & Germer were studying the reflection of electrons from Ni target. The Ni target was subjected to such a heat treatment that the reflection became anomalous.

→ The reflected intensity showed striking maxima & minima.

→ Thus they suspected that the electrons are diffracted like X-rays!

→ The D & G - Arrangement is shown in fig.

Apparatus -

- (1) Electron Gun (A)
- (2) Tungsten filament (F)
- (3) Electron detector (C)
- (4) Circular Scale (S) → angle  $[29^\circ \text{ to } 90^\circ]$

The Davisson and Germer experimental arrangement is shown in fig. The apparatus consists of a cathode ray tube where the  $e^-$ s are produced and obtained in an electronic beam of known frequency.

→ The cathode ray tube consists of a tungsten filament (F) where the  $e^-$ s are accelerated in the field. After  $e^-$ s transmitted from suitable slits they form a fine beam & directly fall on to Ni target.

→ The  $e^-$ s acting as waves, diffracted in different directions. The distribution is measured by an  $e^-$  detector which is connected to galvanometer.

→ The  $e^-$  detector [Faraday cylinder (C)] can move on a circular graduated scale (S) b/w  $29^\circ$  &  $60^\circ$  to receive & reflect  $e^-$ s.

→  $e^-$  detector consists of two walls which are insulated from each other.

→

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A retarding potential is maintained between them so that only fast moving  $e^{-}$ s coming from gun may enter it. [electron detector & walls]

In this way the galvanometer deflection is only due to electrons coming from electron Gun.

→ The accelerating potential  $V$  is given a low value & crystal is set at any arbitrary angle  $\theta$ .

The electron detector is moved to various positions on scale & Galvanometer current is noted for each position.

Hence it should be remembered that galvanometer current is a measurement of intensity of diffracted beam.

A graph is plotted between Galvanometer current and angle  $\theta$  between incident & diffracted beam it comes as shown in

teo  
to  
with

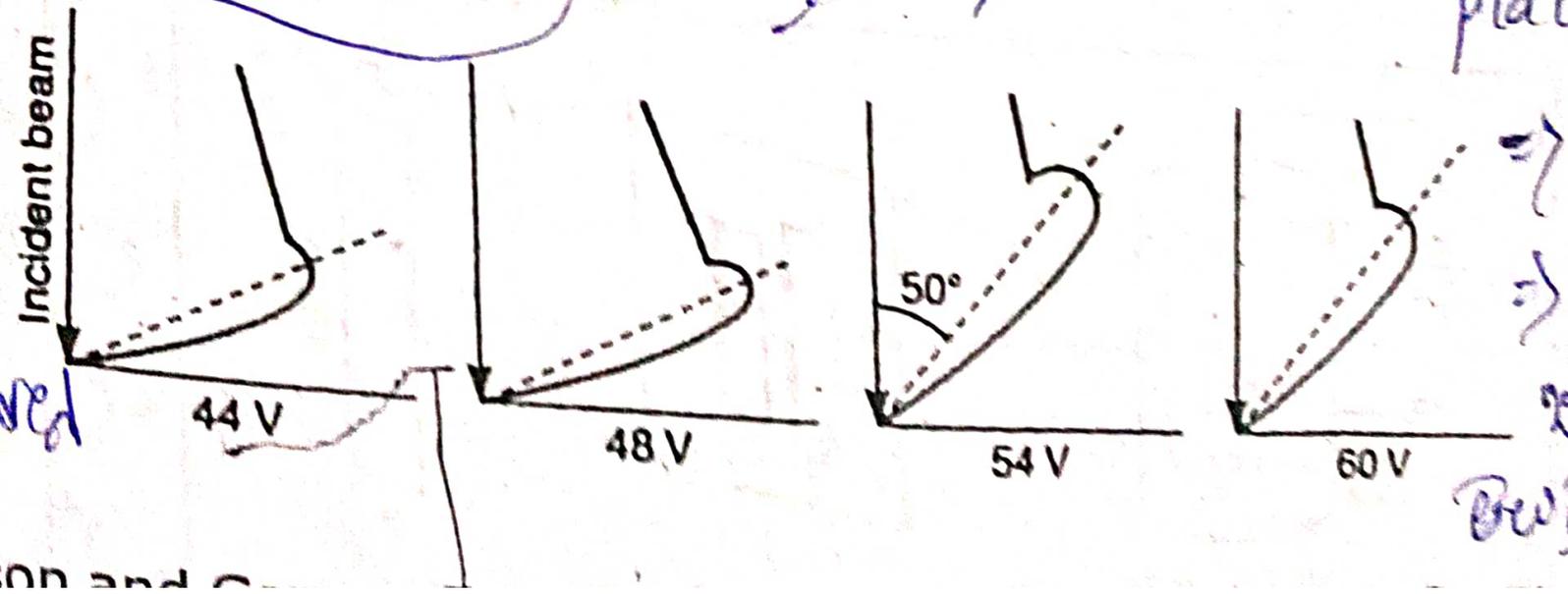
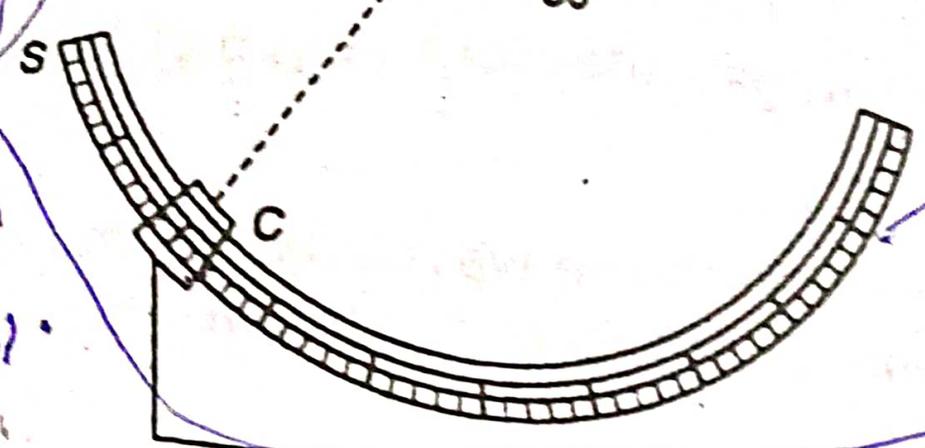
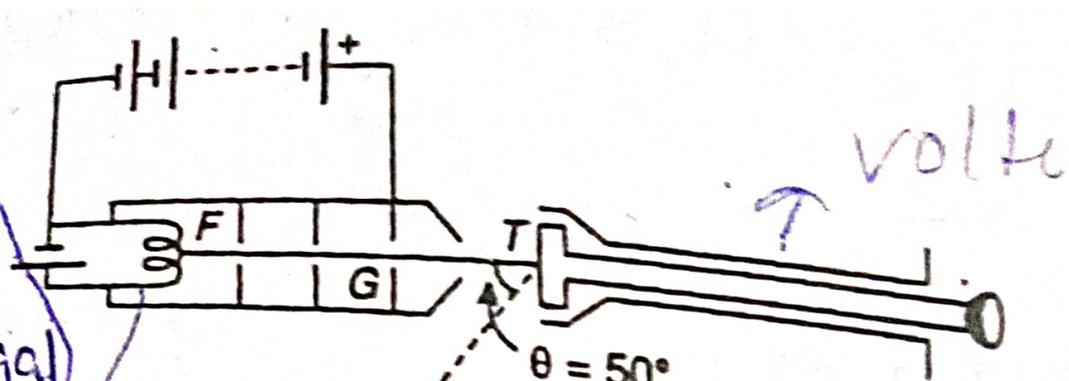
Potential

Plane → red  
6.2/100.

ve electrons

29 batteries 20V

fine sharp beam end 60V  
2000V → 200V



60V

60V

It is observed from graph that  
are

- (i) with ↑ potential bump moves ↑ [UW]  
(ii) the bump becomes most prominent  
in curve for (54V)  $e^-$ s at  $50^\circ$   
(iii) At higher potentials, bump gradually  
disappears.

⇒ According to De-Broglie, wavelength  
associated with  $e^-$  accelerated  
through a potential 'V' is given by.

$$\lambda = \frac{12.26}{\sqrt{V}} \text{ \AA}$$

⇒ wavelength associated with 54V  $e^-$ s

$$\lambda = \frac{12.26}{\sqrt{54}} \Rightarrow \boxed{\lambda = 1.67 \text{ \AA}}$$

In case of x-ray diffraction,  
the wavelength of x-rays  
According to Bragg's eqn 5]  
is  $\lambda = 1.65 \text{ \AA}$

This is almost equal to  $\lambda$  computed  
from De-Broglie hypothesis.

Hence D&G exper confirm Debroglie  
cmw.

# Photo Electric Effect

Experimentally

PEE :- "when a beam of light of suitable frequency (or) wavelength, is incident on a metal surface then the electrons are emitted from that metal surface"

⇒ This phenomenon is known as PEE.  
The emitted electrons are called as photo electrons

⇒ The current created by photo  $e^-$  is called photo current  
Experimental Study of photoelectric Effect

The Experimental Arrangement consists of the following apparatus.  
Mainly 2 parts

- 1) Quartz Bulb
- 2) External circuit

↓  
The small metal plate 'A' and an electrode 'B' is placed inside the evacuated quartz bulb.

The plate 'A' is connected to the negative terminal of battery while 'B' is connected to the +ve terminal of battery through Galvanometer G.

Process

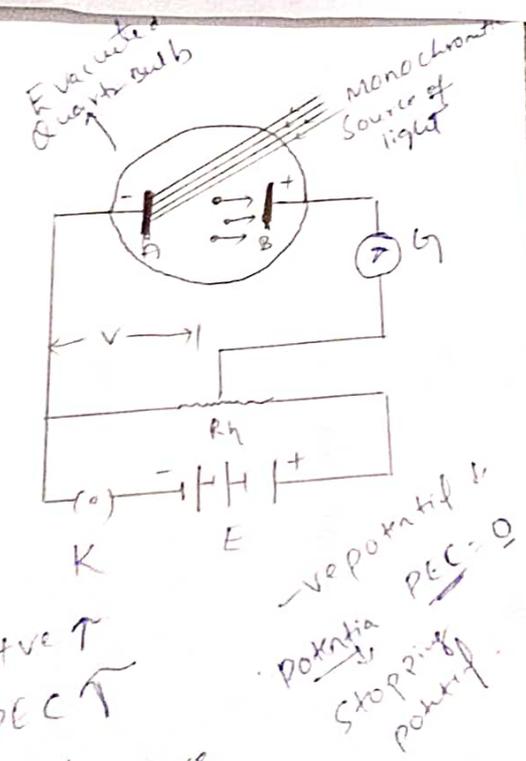
In the absence of light, no electrons flow in the circuit and hence there is no deflection in Galvanometer.

But when monochromatic light is allowed to fall on plate A, electrons are emitted and travel towards electrode B and hence current starts flowing in circuit which is indicated by Galvanometer. This current is known as photo current.

If the potential of (B) made -ve

⇒ photo electric current does not immediately drop to zero because the photo electrons are emitted from (A) with finite velocity.

→ If -ve potential further increased the photo current decreases and finally becomes zero at a particular value known as stopping potential, it is denoted by  $(V_0)$



+ve T  
PEC ↑

∴ The -ve potential of electrode at which PEC is zero is called stopping potential  $(V_0)$

## ① De-Broglie's hypothesis or [DB $\lambda$ ]

→ According to De Broglie hypothesis

moving particle is associated with a wave known as De-Broglie wave.

The wavelength of matter wave is given by

$$\lambda = \frac{h}{mv} = \frac{h}{p}$$

5 things to Remember.

$$\left. \begin{array}{l} \text{① Planck's } E = h\nu \\ \text{② Einstein's } E = mc^2 \end{array} \right\} \lambda = \frac{h}{mc} \Rightarrow C = V$$

③ KE of  $e^-$

④ KE of charge

⑤  $e^{-nS} \Rightarrow E = KE = \frac{1}{2} m_0 v^2 = eV$

- According to Planck's theory  
energy of photon,

$$E = h\nu \quad \text{--- (1)}$$

According to eqn of Einsteins  
mass - Energy radiation

$$E = mc^2 \quad \text{--- (2)}$$

from (1) & (2)  $h\nu = mc^2$

$$\nu = \frac{mc^2}{h} \Rightarrow \frac{c}{\lambda} = \frac{mc^2}{h}$$

$$\lambda = \frac{h}{mc}$$

In case of material particle with velocity  
 $v$ .

$$\lambda = \frac{h}{mv}$$

(2) KE of particle is  $E \Rightarrow E = \frac{1}{2}mv^2$

$$mv^2 = \frac{2E}{1} \Rightarrow mv^2 = 2mE$$

$$p^2 = 2mE$$

$$p = \sqrt{2mE}$$

$$\Rightarrow \lambda = \frac{h}{\sqrt{2mE}}$$

$\vec{p} \rightarrow \vec{p} \rightarrow$  sub  
p eliminated  
 $E \rightarrow$  new  
comer

(4) when charged particle carrying a charge  $q$  & potential difference  $V$   
 $\Rightarrow KE, E = qV \text{ --- (6)}$

from (5) & (6)

$$\lambda = \frac{h}{\sqrt{2mqV}}$$

$E \rightarrow$  Sub.  
 $E \rightarrow$  eliminated  
 $q, V \rightarrow$  new come

(5) In case of electrons.

$$E = KE = \frac{1}{2} m_0 v^2 = eV$$

$$v = \frac{\sqrt{2eV}}{m_0}$$

$$\Rightarrow \lambda = \frac{h}{m_0 v} \Rightarrow \lambda = \frac{h}{m_0 \frac{\sqrt{2eV}}{m_0}} = \frac{h}{\sqrt{2eV}}$$

$$\lambda = \frac{h}{\sqrt{2eV m_0}}$$

$$\lambda = \frac{12.26}{\sqrt{V}} \text{ \AA}$$

$$\therefore V = 100V \Rightarrow \lambda = \frac{12.26}{\sqrt{100}} \text{ \AA} = 1.226 \text{ \AA}$$

# Heisenberg uncertainty principle

It is impossible to specify/measure precisely and simultaneously the values of both members of particular physical variables that describes the behaviour of an atomic system.

- (1) position-momentum uncertainty
- (2) Energy-time uncertainty principle

→ The order of magnitude of uncertainties in the knowledge of 2 variables must be at least planks constant  $h$ .

→ Considering the pair of physical variables as 'a' and 'b' then  $\Delta a \cdot \Delta b \cong h$ .

① position and momentum uncertainty:-

$\Delta x$  = uncertainty in measurement of position  
 $\Delta p$  = " " " " momentum

⇒ According to principle

$\Delta x \cdot \Delta p \cong h$  and also → They won't less than order of  $\frac{h}{4\pi}$   $\Delta x \cdot \Delta p \geq \frac{h}{4\pi}$

## (11) Energy-time uncertainty

$$\Delta E = \text{UIM} \text{ of Energy}$$

$$\Delta p_x = \text{UIM} \text{ of momentum}$$

$$\Delta t = \text{UIM} \text{ of time}$$

$$\Rightarrow \Delta E \cdot \Delta t \cong h \text{ and } \Delta E \cdot \Delta t \geq \frac{h}{4\pi}$$

Proof:  $\therefore$  of  $\Delta E \Delta t \geq \frac{h}{4\pi}$

Any particle mass  $m_0$  & velocity  $v_n$   
 $\Rightarrow$  KE

$$E = \frac{1}{2} m_0 v_n^2 = \frac{m_0 v_n^2}{2 m_0} = \frac{p_x}{2 m_0}$$

$$\Delta E = \frac{2 \cdot p_x \Delta p_x}{2 m_0}$$

$$\Delta p_x \cdot p_x = m_0 \Delta E$$

$$\Delta p_x = \frac{m_0 \Delta E}{p_x} = \frac{m_0 \Delta E}{m_0 v_n} = \frac{\Delta E}{v_n}$$

$$\boxed{\Delta E = v_n \Delta p_x} \Rightarrow \Delta p_x = \frac{\Delta E}{v_n}$$

We know  $v_n = \frac{\Delta n}{t} \Rightarrow v_n \Delta t = \Delta n$

$$\boxed{\Delta t = \frac{\Delta n}{v_n}} \Rightarrow \Delta n = \Delta t \cdot v_n$$

$$\Delta E \cdot \Delta t = v_n \Delta p_x \frac{\Delta n}{v_n} = \Delta n \cdot \Delta p_x \cong h$$

$$\Rightarrow \text{hence } \Delta E \cdot \Delta t \cong h \quad \text{or} \quad \boxed{\Delta E \cdot \Delta t \geq \frac{h}{4\pi}} \text{ from (1)}$$

# SCHRODINGER Wave

## Equation

- (1) Time dependent SWE
- (2) Time Independent SWE

## SW Equations:

To study the atomic systems, we need Equation of motion.

Similarly to study de Broglie's waves associated with particles and describes the motion of particles, we need a fundamental equations in Quantum mechanics known as Schrodinger wave Equations.

Schrodinger presented the wave-wave Equations as a development of De-Broglie ideas of wave properties of matter for this process purpose, Schrodinger introduced a mathematical function ( $\psi$ ) known

as wave function which is the variable quantity associated with the moving particle and is also complex function of the space co-ordinates of the particle and time.

$$\text{Let, } \psi = A e^{i(kx - \omega t)} \quad \text{--- (1)}$$

$$\frac{\partial \psi}{\partial x} = ik A e^{i(kx - \omega t)}$$

$$\frac{\partial^2 \psi}{\partial x^2} = A \cdot i^2 k^2 e^{i(kx - \omega t)}$$

$$\frac{\partial^2 \psi}{\partial x^2} = -k^2 \psi \quad \text{--- (2)}$$

But we know that  $k = \frac{2\pi}{\lambda}$  &  $\lambda = \frac{h}{p}$

$$\therefore k = \frac{2\pi \cdot p}{h}$$

$$k^2 = \frac{4\pi^2 p^2}{h^2}$$

$$\text{from (2)} \quad \frac{\partial^2 \psi}{\partial x^2} = -\frac{4\pi^2 p^2}{h^2} \psi$$

$$p^2 \psi = -\frac{h^2}{4\pi^2} \frac{\partial^2 \psi}{\partial x^2} \quad \text{--- (3)}$$

$$\textcircled{1} \rightarrow \psi = A e^{i(kx - \omega t)} \quad \textcircled{1}$$

Differentiating  $\textcircled{1}$  w.r to  
time then

$$\frac{\partial \psi}{\partial t} = +A e^{i(kx - \omega t)} [-i\omega]$$

$$\frac{\partial \psi}{\partial t} = -i\omega \psi \quad \textcircled{4}$$

we know that

$$E = h\nu = \frac{h\omega}{2\pi}$$

$$\omega = \frac{2\pi E}{h}$$

from  $\textcircled{4}$  we get

$$\frac{\partial \psi}{\partial t} = -i \frac{2\pi E}{h} \psi$$

$$\textcircled{2} \quad \frac{\partial \psi}{\partial t} \cdot \frac{h}{2\pi i} = -E\psi$$

$$E\psi = -\frac{\partial \psi}{\partial t} \cdot i\hbar \quad \left[ \because \frac{h}{2\pi} = \hbar \right]$$

$$E\psi = i\hbar \frac{\partial \psi}{\partial t} \quad \textcircled{5}$$

But, we know that

$$E = KE + PE$$

$$E = \frac{p^2}{2m} + V$$

$$E\psi = \frac{p^2}{2m}\psi + V\psi \quad (6)$$

Sub eqns (3) & (5) in (6)  
we get

$$i\hbar \frac{\partial \psi}{\partial t} = -\frac{\hbar^2}{4\pi^2} \frac{\partial^2 \psi}{\partial x^2} \times \frac{1}{2m} + V\psi$$

$$i\hbar \frac{\partial \psi}{\partial t} = -\frac{\hbar^2}{2m} \frac{\partial^2 \psi}{\partial x^2} + V\psi$$

$$i\hbar \frac{\partial \psi}{\partial t} = \left[ \frac{-\hbar^2 \Delta^2 + V}{2m} \right] \psi \rightarrow (7)$$

[in case of 3-dimensional]

$$\hat{E}\psi = \hat{H}\psi \quad \text{where}$$

$$\hat{E} = i\hbar \frac{\partial}{\partial t} = \text{Energy operator}$$

$$\hat{H} = \frac{-\hbar^2}{2m} \Delta^2 + V = \text{Hamiltonian operator}$$

# Time Independent Schrödinger

## Wave equation

In many situations, PE do not depend on time. In those,

The forces that are acting on the PE ( $V$ ) depend on the position of the particle. Under these conditions,

$\psi(x,t)$  can be written as the product of a function of  $x$ -alone and the other the functions of  $t$  alone.

$$\text{Let, } \psi(x,t) = A e^{i(kx - \omega t)}$$

$$\psi(x,t) = A e^{pkx} e^{-i\omega t}$$

$$\psi(x,t) = \phi(x) e^{-i\omega t}$$

$$\frac{\partial \psi}{\partial t} = -i\omega \phi(x) e^{-i\omega t}$$

$$\frac{\partial \psi}{\partial t} = -i\omega \psi$$

$$\frac{\partial \psi}{\partial t} = -i \frac{2\pi E}{h} \psi \quad \left[ \omega = \frac{2\pi E}{h} \right]$$

(2)

$$\Rightarrow E\psi = p \frac{h}{2\pi} \frac{\partial \psi}{\partial t}$$

$$E\psi = i \frac{h}{2\pi} \frac{\partial \psi}{\partial t}$$

$$\boxed{E = i\hbar \frac{\partial}{\partial t}}$$

$$E\psi = i\hbar \frac{\partial \psi}{\partial t} \Rightarrow$$

$$E\psi = \underline{E\psi}$$

on differentiating equation (1) w.r.t  $x$ , then

$$\frac{\partial \psi}{\partial x} = e^{-i\omega t} \frac{\partial \phi}{\partial x}$$

$$\frac{\partial^2 \psi}{\partial x^2} = e^{-i\omega t} \frac{\partial^2 \phi}{\partial x^2}$$

According to Schrodinger, Time dependent equation,

$$E\psi = -\frac{\hbar^2}{2m} \frac{\partial^2 \psi}{\partial x^2} + V\psi \quad \text{--- (4)}$$

Sub (1) & (3) in eqn (4) we get.

$$E \cdot \phi(x) e^{-i\omega t} = -\frac{\hbar^2}{2m} e^{-i\omega t} \frac{\partial^2 \phi}{\partial x^2} + V\phi(x) e^{-i\omega t}$$

$$E\phi(x) = -\frac{\hbar^2}{2m} \frac{\partial^2 \phi}{\partial x^2} + V\phi(x)$$

$$E\phi(x) - V\phi(x) = -\frac{\hbar^2}{2m} \frac{\partial^2 \phi}{\partial x^2}$$

$$\frac{\hbar^2}{2m} \frac{\partial^2 \phi}{\partial x^2} + (E - V)\phi(x) = 0$$

$$\frac{\partial^2 \phi}{\partial x^2} + \frac{2m}{\hbar^2} (E - V)\phi(x) = 0$$

In General,  $\frac{\partial^2 \psi}{\partial x^2} + \frac{2m}{\hbar^2} (E - V)\psi = 0$ .

In case of 3-Dimensions

$$\left[ \nabla^2 \psi + \frac{2m}{\hbar^2} (E - V)\psi = 0 \right] \text{--- (5)}$$

for a free particle  $V=0$ .

∴ from (5)

$$\nabla^2 \psi + \frac{2mE}{\hbar^2} \psi = 0$$

eqn (5)

$$\nabla^2 \psi = -\frac{2m}{\hbar^2} (E - V)\psi$$

$$-\frac{\hbar^2}{2m} \nabla^2 \psi = E\psi - V\psi$$

$$\left[ \frac{\hbar^2}{2m} \nabla^2 + V \right] \psi = E\psi$$

$$\hat{H} \psi = \hat{E} \psi$$

Momentum of System along x-axis  
 Before collision =  $p_p + p_e = \frac{h\nu}{c} + 0$   
 After collision =  $p_p' + p_e' = \frac{h\nu' \cos \theta}{c} + m\nu \cos \phi$   
 According to law of conservation of momentum along x-axis  
 $\frac{h\nu}{c} = \frac{h\nu' \cos \theta}{c} + m\nu \cos \phi$  (1)  
 $h\nu = h\nu' \cos \theta + m\nu c \cos \phi$   
 $m\nu c \cos \phi = h\nu - h\nu' \cos \theta$  (2)

System along y-axis  
 Before collision = 0  
 After collision =  $p_p' + p_e' = \frac{h\nu' \sin \theta}{c} + m\nu \sin \phi$   
 According to law of conservation of momentum along y-axis  
 $0 = \frac{h\nu'}{c} \sin \theta + m\nu \sin \phi$   
 $0 = h\nu' \sin \theta + m\nu c \sin \phi$   
 $h\nu' \sin \theta = -m\nu c \sin \phi$  (3)  
 $m\nu c \sin \phi = -h\nu' \sin \theta$  (3)

Total energy of System before collision =  $E_p + E_e = h\nu + mc^2$   
 After collision =  $E_p' + E_e' = h\nu' + mc^2$   
 According to law of conservation of Energy:  
 $E_p + E_e = E_p' + E_e'$   
 $h\nu + mc^2 = h\nu' + mc^2$   
 $mc^2 = h(\nu - \nu') + mc^2$  (S.N.S.)  
 $m^2 c^4 = h^2 (\nu - \nu')^2 + m^2 c^4 + 2h(\nu - \nu') mc^2$   
 $m^2 c^4 = h^2 \nu^2 + h^2 \nu'^2 - 2h\nu\nu' + m^2 c^4 + 2h\nu mc^2 - 2h\nu' mc^2$  (5)

(1) + (3)  
 $LHS + LHS = RHS + RHS$   
 $m^2 \nu^2 c^2 \cos^2 \phi + m^2 \nu^2 c^2 \sin^2 \phi = h^2 \nu^2 + h^2 \nu'^2 \cos^2 \theta - 2h\nu h\nu' \cos \theta + h^2 \nu'^2 \sin^2 \theta + 2h\nu mc^2 - 2h\nu' mc^2$   
 $m^2 \nu^2 c^2 (\cos^2 \phi + \sin^2 \phi) = h^2 \nu^2 - 2h^2 \nu \nu' \cos \theta + h^2 \nu'^2 (\cos^2 \theta + \sin^2 \theta) + 2h\nu mc^2 - 2h\nu' mc^2$   
 $m^2 \nu^2 c^2 = h^2 \nu^2 + h^2 \nu'^2 - 2h^2 \nu \nu' \cos \theta + 2h\nu mc^2 - 2h\nu' mc^2$  (4)

$m^2 c^4 - m^2 \nu^2 c^2 = h^2 \nu^2 + h^2 \nu'^2 - 2h^2 \nu \nu' \cos \theta + 2h\nu mc^2 - 2h\nu' mc^2 - h^2 \nu^2 - h^2 \nu'^2 + 2h^2 \nu \nu' \cos \theta$   
 $m^2 c^2 [c^2 - \nu^2] = -2h^2 \nu \nu' [1 - \cos \theta] + 2h(\nu - \nu') mc^2 + m^2 c^4$   
 $\frac{m^2 c^2}{c^2 - \nu^2} (c^2 - \nu^2) = \frac{m^2 c^2}{c^2 - \nu^2} c^2 (c^2 - \nu^2) = -2h^2 \nu \nu' (1 - \cos \theta) + 2h(\nu - \nu') mc^2 + m^2 c^4$   
 $m^2 c^4 = -2h^2 \nu \nu' (1 - \cos \theta) + 2h(\nu - \nu') mc^2 + m^2 c^4$   
 $2h^2 \nu \nu' (1 - \cos \theta) = 2h(\nu - \nu') mc^2$

$$2h(v-v')mc^2 = 2h^2vv'(1-\cos\theta)$$

$$\frac{v-v'}{vv'} = \frac{h}{mc^2} (1-\cos\theta)$$

$$\frac{1}{v'} - \frac{1}{v} = \frac{h}{mc^2} (1-\cos\theta)$$

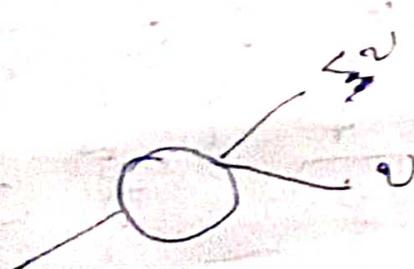
$$\frac{c}{v'} - \frac{c}{v} = \frac{h}{mc} (1-\cos\theta)$$

$$\frac{1}{\lambda'} - \frac{1}{\lambda} = \Delta\lambda = \frac{h}{mc} (1-\cos\theta)$$

$$\Delta\lambda = \frac{h}{mc} \left(1 - \left(1 - 2\sin^2\frac{\theta}{2}\right)\right)$$

$$\Delta\lambda = \frac{2h}{mc} \sin^2\frac{\theta}{2}$$

$\Delta\lambda$  = change in wavelength



# Quantum Mechanics

Paper DC

Date \_\_\_\_\_

Page \_\_\_\_\_

3 parts

- part - 1
- (1) PEE & Compton effect (3)  $10M$
  - (2) Einstein PEE  $10M$
  - (3) Quantum mechanics postulates  $10M$

part - 2 :-

- last w
- (1) Debroglie wavelength
  - (2) Heisenberg uncertainty principle
  - (3) matter waves properties  $10M$
  - (4) LA phase velocity & group v difference  $30M$
- last w
- (5) Davison germ experiment
  - (6) Gamma ray microscope  $10M$

part - 3 :- w, Schrodinger time d  
last w

- 3) LA properties of wave function -
- 4) LA eigen values eigen functions -